

Metals Magnetic systems

Ralph Gebauer

Cape Town, July 2008

Metallic systems: k-points and smearing:

Let us consider iron, in the bcc phase.

In the first part of this exercise, we **neglect magnetism** and perform non-spin polarized calculations.

Task 1: Create a pwscf input file for scf total energy calculations in this system with the following features:

- Use the **Fe.pz-nd-rrkjus.UPF** pseudopotential
- Cutoffs: **30 Ry** (wavefunctions), **250 Ry** (charge density)
- Lattice constant: **5.38 (a.u.)**
- K-point mesh: **7×7×7** with shift **1×1×1**
- **Marzari-Vanderbilt** smearing with **degauss=0.03 Ry**

Solution 1:

```
&control
  calculation = 'scf',
  prefix = 'Fe_bcc',
  pseudo_dir = './',
  outdir = './'
/
&system
 ibrav = 3,
  celldm(1) = 5.38,
  nat = 1,
  ntyp = 1,
  nbnd = 10
  ecutwfc = 30.,
  ecutrho = 250.,
  occupations='smearing'
  smearing='marzari-vanderbilt'
  degauss=0.03
/
&electrons
  mixing_beta = 0.7
/
ATOMIC_SPECIES
Fe 55.845 Fe.pz-nd-rrkjus.UPF
ATOMIC_POSITIONS (alat)
Fe 0.0 0.0 0.0
K_POINTS (automatic)
7 7 7 1 1 1
```

bcc-lattice
lattice constant

increased number of bands
cutoffs

metallic smearing

Metallic systems: k-points and smearing:

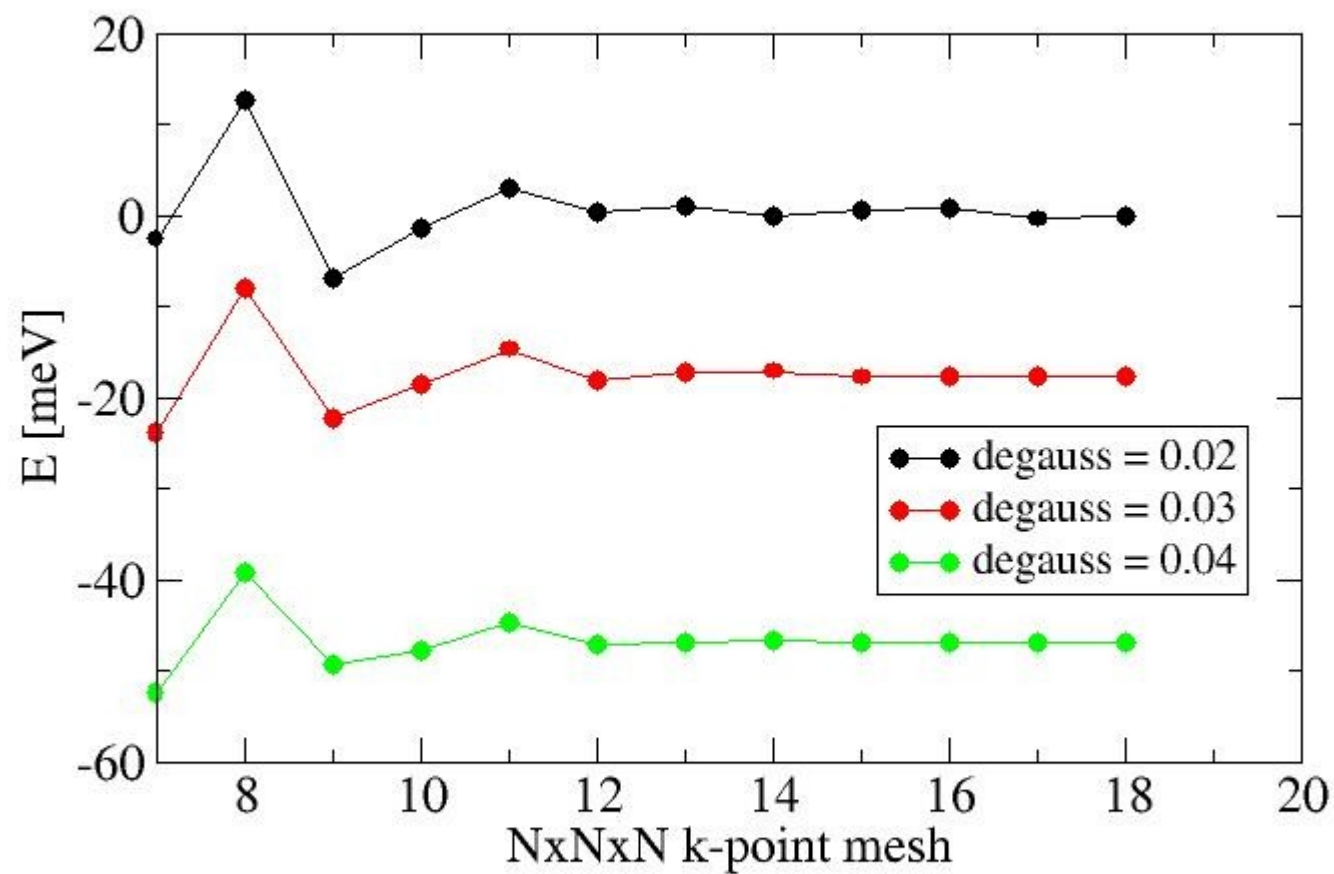
Task 2: Compute the total energy for this system with the following values for degauss and k-point meshes:

degauss = 0.02, 0.03, 0.04

k-points: $N \times N \times N$ ($1 \times 1 \times 1$) with $N=7,8,9, \dots, 18$

Plot the convergence of the total energy with respect to the number of k-points for the different smearing values.

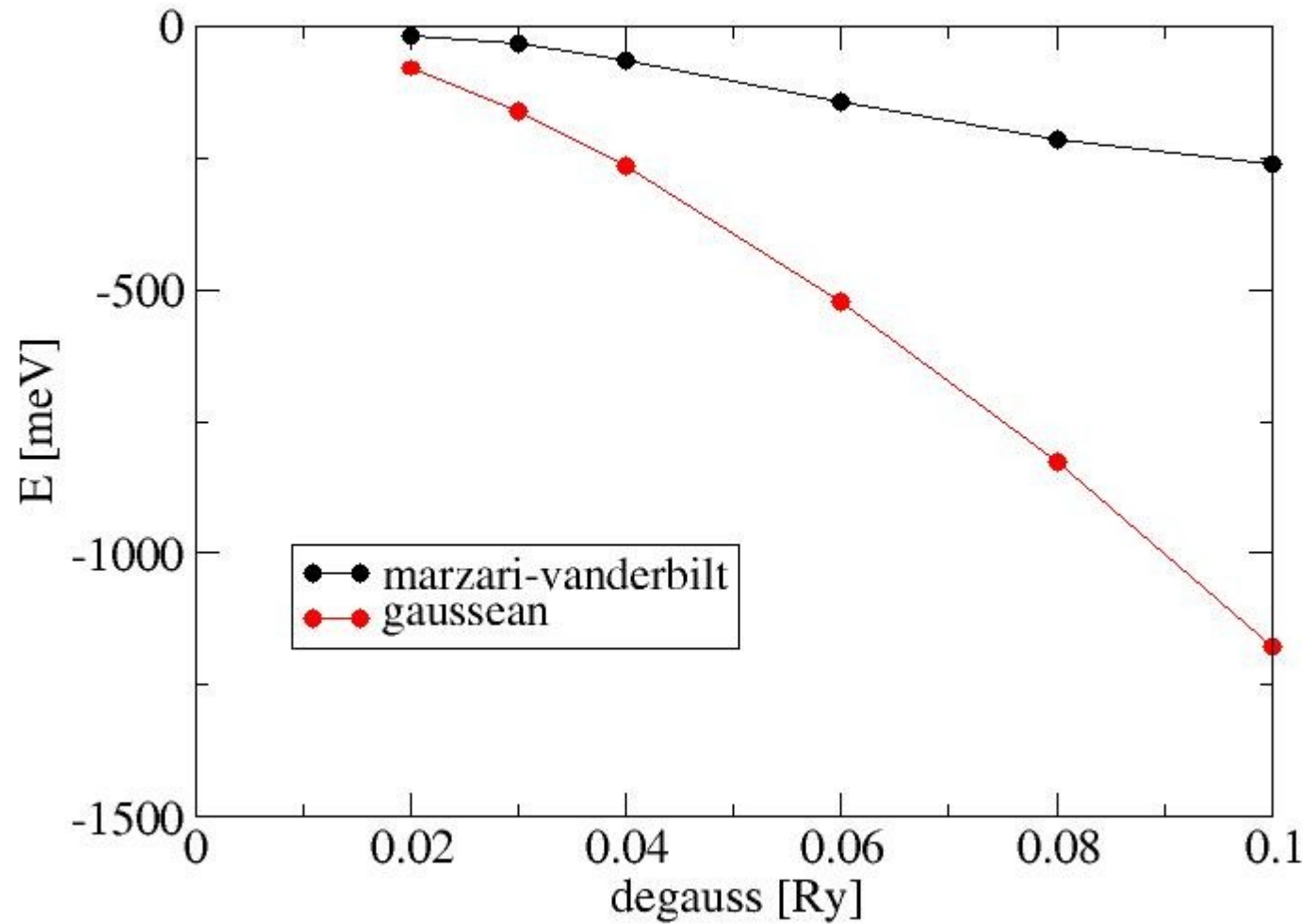
Solution2:



Metallic systems: k-points and smearing:

Task 3: Choose a $14 \times 14 \times 14$ ($1 \times 1 \times 1$) k-point mesh and plot the total energy as a function of degauss for Marzari-Vanderbilt and Gaussian smearing

Solution3:



Magnetism and spin polarized systems

So far in this course: we “neglected” spin of the electrons:

$$n(\mathbf{r}) = 2 \sum_{i=1}^{N_{el}/2} |\psi_i(\mathbf{r})|^2$$

Implicitly, we have assumed that spin-up and spin-down wavefunctions are the same:

$$\psi_i^{\uparrow}(\mathbf{r}) = \psi_i^{\downarrow}(\mathbf{r})$$

Advantage: we need only $N_{el}/2$ orbitals

Spin polarized systems

For the case of spin-polarized systems, we drop this constraint and
Use N orbitals:

$$\left\{ \psi_i^\uparrow \right\}, \quad i = 1, \frac{N}{2} \quad \left\{ \psi_i^\downarrow \right\}, \quad i = 1, \frac{N}{2}$$

$$n^\uparrow(\mathbf{r}) = \sum_i \left| \psi_i^\uparrow(\mathbf{r}) \right|^2 \quad n^\downarrow(\mathbf{r}) = \sum_i \left| \psi_i^\downarrow(\mathbf{r}) \right|^2$$

$$n(\mathbf{r}) = n^\uparrow + n^\downarrow$$

$$m(\mathbf{r}) = n^\uparrow - n^\downarrow$$

Spin polarized systems

Look at how the total energy depends on the two charge densities:

$$E[n^\uparrow, n^\downarrow] = E_{kin} + E_{ext}[n^\uparrow + n^\downarrow] + E_{Har}[n^\uparrow + n^\downarrow] + E_{xc}[n^\uparrow, n^\downarrow]$$

Calculate the functional derivative in order to obtain the Hamiltonian:

$$\frac{\delta E}{\delta n^\uparrow(\mathbf{r})} \neq \frac{\delta E}{\delta n^\downarrow(\mathbf{r})}$$

Spin polarized systems

We can also write the exchange–correlation energy as $E_{xc}[n, m]$.

Since we have:

$$\frac{\delta n}{\delta n^\uparrow(\mathbf{r})} = \frac{\delta n}{\delta n^\downarrow(\mathbf{r})} = 1$$

$$\frac{\delta m}{\delta n^\uparrow(\mathbf{r})} = 1$$

$$\frac{\delta m}{\delta n^\downarrow(\mathbf{r})} = -1$$

The Hamiltonian looks like:

$$\left[\frac{-\hbar^2}{2m} \nabla^2 + V_{ext} + V_{Har} + V_{xc} + B_{xc} \right] \psi_i^\uparrow = \epsilon_i \psi_i^\uparrow$$

$$\left[\frac{-\hbar^2}{2m} \nabla^2 + V_{ext} + V_{Har} + V_{xc} - B_{xc} \right] \psi_j^\downarrow = \epsilon_j \psi_j^\downarrow$$

Spin polarized systems

$$\begin{aligned}
 & \left[\frac{-\hbar^2}{2m} \nabla^2 + V_{ext} + V_{Har} + V_{xc} + B_{xc} \right] \psi_i^\uparrow = \epsilon_i \psi_i^\uparrow \\
 & \left[\frac{-\hbar^2}{2m} \nabla^2 + V_{ext} + V_{Har} + V_{xc} - B_{xc} \right] \psi_j^\downarrow = \epsilon_j \psi_j^\downarrow
 \end{aligned}$$

Magnetic systems:

We now come back to iron, in the bcc phase.

We now **take into account magnetism** and perform spin polarized calculations.

Task 4: Create a pwscf input file for scf total energy calculations in this system with the following features:

- Use the **Fe.pz-nd-rrkjus.UPF** pseudopotential
- Cutoffs: **30 Ry** (wavefunctions), **250 Ry** (charge density)
- Lattice constant: **5.38 (a.u.)**
- K-point mesh: **14×14×14** with shift **1×1×1**
- **Marzari-Vanderbilt** smearing with **degauss=0.03 Ry**
- **spin polarized**

Solution 4:

```

&control
  calculation = 'scf',
  prefix = 'Fe_bcc',
  pseudo_dir = './',
  outdir = './'
/
&system
 ibrav = 3,
  celldm(1) = 5.38,
  nat = 1,
  ntyp = 1,
  nbnd = 20
  ecutwfc = 30.,
  ecutrho = 250.,
  occupations='smearing'
  smearing='gauss'
  degauss=0.03
  nspin=2
  starting_magnetization(1)=0.2
/
&electrons
  mixing_beta = 0.7
/
ATOMIC_SPECIES
Fe 55.845 Fe.pz-nd-rrkjus.UPF
ATOMIC_POSITIONS (alat)
Fe 0.0 0.0 0.0
K_POINTS (automatic)
14 14 14 1 1 1

```

increased number of bands

set nspin=2

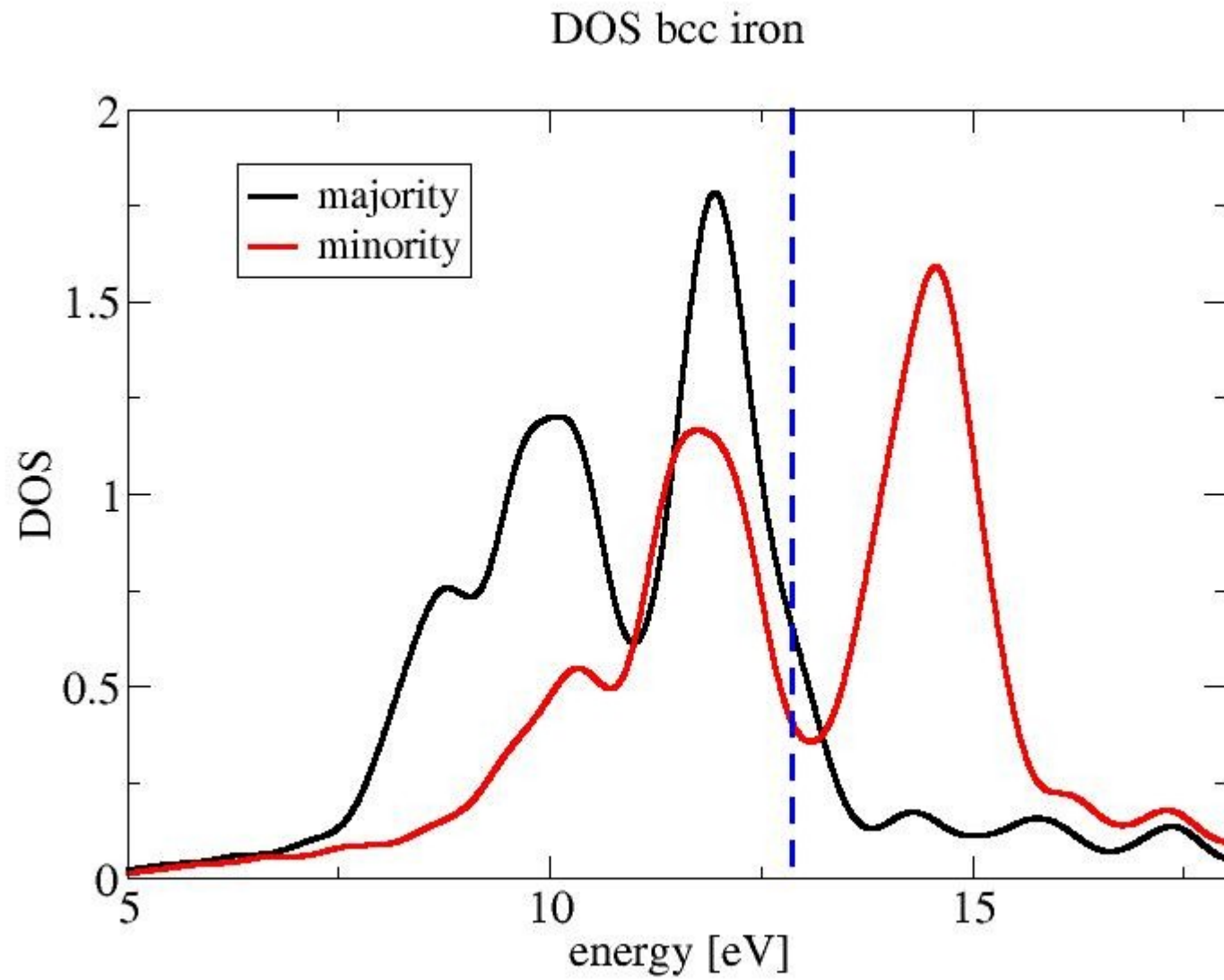
don't forget: starting_magnetization

Magnetic systems: Density of states

Task 5: Using the input file from task 4, calculate the total energy of magnetic bcc iron. Then, use the input file below for the postprocessing program dos.x. Plot the DOS of magnetic iron.

```
&inputpp  
outdir='./'  
prefix='Fe_bcc'  
fildos='Fe.dos'  
Emin=5., Emax=18., DeltaE=0.005  
/
```

Solution 5:



Magnetic systems: Density of states

Task 6: Do a non-selfconsistent calculation of magnetic bcc iron using the “tetrahedron method” for determining the occupations. Then calculate the DOS again (using dos.x). Plot the DOS.

Solution 6:

